# Remarks on the Transformations of Space and Time 

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Space and time transformations from a "stationary" isotropic inertial system $S_{0}$ to any other inertial system S are shown to imply complete physical equivalence between the three possible pairs of assumptions chosen among the following: A1. Lorentz contraction of bodies moving with respect to $S_{0} ;$ A2. Larmor retardation of clocks moving with respect to $S_{0} ; A 3$. Two-way velocity of light equal to $c$ in all inertial systems and in all directions. The empirical evidence supporting $A 2$ and A3 is therefore in favour of A1 as well.

## 1. Introduction

We suppose that a ("stationary") inertial reference frame $S_{0}$ exists in which Maxwell's equations hold. A well known consequence of these equations is that the velocity of light in the vacuum is $c$ in all directions. Therefore in $S_{0}$ clocks must be synchronised by using Einstein's procedure [1]. Space and time variables $x_{0}, y_{0}, z_{0}, t_{0}$ are thus available in all points of $S_{0}$ and the velocities of moving bodies (and of other inertial systems) can be measured. The most general form of space-time transformation from $S_{0}$ to a different inertial system $S(x, y, z, t)$ is:

$$
\left\{\begin{align*}
x & =f\left(x_{0}, y_{0}, z_{0}, t_{0}\right)  \tag{1}\\
y & =g\left(x_{0}, y_{0}, z_{0}, t_{0}\right) \\
z & =h\left(x_{0}, y_{0}, z_{0}, t_{0}\right) \\
t & =e\left(x_{0}, y_{0}, z_{0}, t_{0}\right)
\end{align*}\right.
$$

where $f, g, h, e$ are four functions of the $S_{0}$ space coordinates $x_{0}, y_{0}, z_{0}$, and time $t_{0}$. In reducing the generality of (1) one can follow two guiding principles:
A. Empty space is homogeneous, that is all points have the same physical properties, and is isotropic, that is all directions are physically equivalent.
B. Time is homogeneous, that is all properties of space remain the same with passing time.

In Ref. [2] it was shown that, given $A$ and $B$, the transformations (1) must be linear in all variables:

$$
\left\{\begin{array}{l}
x=f_{1} x_{0}+f_{2} y_{0}+f_{3} z_{0}+f_{4} t_{0}+f_{5}  \tag{2}\\
y=g_{1} x_{0}+g_{2} y_{0}+g_{3} z_{0}+g_{4} t_{0}+g_{5} \\
z=h_{1} x_{0}+h_{2} y_{0}+h_{3} z_{0}+h_{4} t_{0}+h_{5} \\
t=e_{1} x_{0}+e_{2} y_{0}+e_{3} z_{0}+e_{4} t_{0}+e_{5}
\end{array}\right.
$$

The twenty coefficients appearing in (2) are constant with respect to $x_{0} y_{0}, z_{0}, t_{0}$ but can naturally depend on the velocity $v$ of $S$ relative to $S_{0}$. One can choose the coordinate systems in $S$ and in $S_{0}$ in such a way that the observer in $S$ sees his origin $(x=y=z=0)$ coincident
with that of $S_{0}$ at $t=0$, and vice versa that the observer in $S_{0}$ sees his origin $\left(x_{0}=y_{0}=z_{0}=0\right)$ coincident with that of $S$ at time $t_{0}=0$. This is equivalent to a suitable choice of the origin of the space co-ordinates and of time in $S$. Symbolically one can write:

$$
\left[x_{0}=y_{0}=z_{0}=t_{0}=0\right] \quad \Rightarrow \quad[x=y=z=t=0]
$$

Inserted in (2) this gives:

$$
\begin{equation*}
f_{5}=g_{5}=h_{5}=e_{5}=0 \tag{3}
\end{equation*}
$$

Assume next that plane ( $x_{0}, y_{0}$ ) coincides with plane ( $x, y$ ) at all times $t_{0}$ :

$$
\left[z_{0}=0\right] \Rightarrow[z=0]
$$

The third of (2) gives:

$$
\begin{equation*}
h_{1}=h_{2}=h_{4}=0 \tag{4}
\end{equation*}
$$

Assume also that plane $\left(x_{0}, z_{0}\right)$ coincides with plane $(x$, $z$ )for all times $t_{0}$ :

$$
\left[y_{0}=0\right] \Rightarrow[y=0]
$$

The second of (2) gives:

$$
\begin{equation*}
g_{1}=g_{3}=g_{4}=0 \tag{5}
\end{equation*}
$$

Finally assume that at time $t_{0}=0$ plane $\left(y_{0}, z_{0}\right)$ coincides with plane $(y, z)$. This is like saying that the relative velocity is parallel to the $x_{0}$ axis [and then perpendicular to plane ( $y_{0}, z_{0}$ ), if the Cartesian co-ordinates are orthogonal]. The condition is then:

$$
\left[t_{0}=x_{0}=0\right] \Rightarrow[x=0]
$$

Given (3), from the first of (2) it follows:

$$
\begin{equation*}
f_{2}=f_{3}=0 \tag{6}
\end{equation*}
$$

Consider next to relative velocity condition by assuming that the origin of $S(x=0)$ seen from $S_{0}$ satisfies the equation $x_{0}=v t_{0}$. By substituting $x=0$ and $x_{0}=v t_{0}$ in the first Eq. (2) and taking into account (3) and (6) it follows

$$
\begin{equation*}
f_{4}=-f_{1} v \tag{7}
\end{equation*}
$$

One can now rewrite transformations (2) using (3)-(7):

$$
\left\{\begin{array}{l}
x=f_{1}\left(x_{0}-v t_{0}\right)  \tag{8}\\
y=g_{2} y_{0} \\
z=h_{3} z_{0} \\
t=e_{1} x_{0}+e_{2} y_{0}+e_{3} z_{0}+e_{4} t_{0}
\end{array}\right.
$$

There is a complete equivalence of the axes $y_{0}$ and $z_{0}$, since the relative velocity is parallel to the $x_{0}$ axis. These axes are chosen arbitrarily, and nothing physical can distinguish them if space is isotropic. Therefore:

$$
\begin{equation*}
g_{2}=h_{3} ; e_{2}=e_{3} \tag{9}
\end{equation*}
$$

It must furthermore be considered that all points in a plane perpendicular to the $x$-axis in $S$ are physically equivalent because the whole co-ordinate system translates rigidly with a local velocity $\vec{v}$ which is everywhere the same. It follows that one is free to assume

$$
\begin{equation*}
e_{2}=0 \tag{10}
\end{equation*}
$$

Thus one can replace (8) by

$$
\left\{\begin{array}{l}
x=f_{1}\left(x_{0}-v t_{0}\right)  \tag{11}\\
y=g_{2} y_{0} \\
z=g_{2} z_{0} \\
t=e_{1} x_{0}+e_{4} t_{0}
\end{array}\right.
$$

The transformation (11) allows one to calculate the velocity of light relative to $S$ : One must: (i) write it down in terms of space and time intervals $\Delta x, \Delta y, \ldots \Delta t_{0}$; (ii) invert it, expressing as functions of $\Delta x, \ldots \Delta t$; (iii) substitute the new result in the equation

$$
c \Delta t_{0}=\left[\Delta x_{0}^{2}+\Delta y_{0}^{2}+\Delta z_{0}^{2}\right]^{1 / 2}
$$

(iv) introduce polar co-ordinates. As shown in [2], the one-way velocity of light $\tilde{c}_{1}(\theta)$ and the two-way velocity $\tilde{c}_{2}(\theta)$ turn out to be given by:

$$
\begin{align*}
\frac{1}{\tilde{c}_{1}(\theta)} & =\frac{e_{1} c+e_{4} \beta}{c f_{1}\left(1-\beta^{2}\right)} \cos \theta+ \\
& +\frac{e_{4}+e_{1} \beta c}{c}\left[\frac{\cos ^{2} \theta}{f_{1}^{2}\left(1-\beta^{2}\right)^{2}}+\frac{\sin ^{2} \theta}{g_{2}^{2}\left(1-\beta^{2}\right)}\right]^{1 / 2} \tag{12}
\end{align*}
$$

and

$$
\frac{1}{\tilde{c}_{2}(\theta)}=\frac{e_{4}+e_{1} \beta c}{c}\left[\begin{array}{l}
\frac{\cos ^{2} \theta}{f_{1}^{2}\left(1-\beta^{2}\right)^{2}}+  \tag{13}\\
+\frac{\sin ^{2} \theta}{g_{2}^{2}\left(1-\beta^{2}\right)}
\end{array}\right]^{1 / 2}
$$

if $\theta$ is the angle between the propagation direction of light and the "absolute" velocity of $S$ (parallel to the $x$ axis).

## 2. Three basic assumptions

The separate consequences of the following three assumptions will now be deduced:

A1. Lorentz-Fitzgerald contraction. A body at rest in $S$ between the points of co-ordinates $x_{1}$ and $x_{2}$ when
seen from $S_{0}$ appears contracted along the $x$ direction according to the equation

$$
\begin{equation*}
x_{02}-x_{01}=\sqrt{1-\beta^{2}}\left(x_{2}-x_{1}\right) \tag{14}
\end{equation*}
$$

A body at rest in $S$ between the points of coordinates $y_{1}$ and $y_{2}$ does not appear contracted along the $y$ direction, that is it satisfies

$$
\begin{equation*}
y_{02}-y_{01}=y_{2}-y_{1} \tag{15}
\end{equation*}
$$

A2. Larmor retardation. A clock at rest in any point of $S$, when seen from $S_{0}$ appears retarded according to the equation

$$
\begin{equation*}
t_{02}-t_{01}=\frac{1}{\sqrt{1-\beta^{2}}}\left(t_{2}-t_{1}\right) \tag{16}
\end{equation*}
$$

A3. Invariance of two-way velocity of light. A flash of light propagating forth and back on any segment $A B$ at rest in $S$ does so with a two-way velocity

$$
\begin{equation*}
\tilde{c}_{2}(\theta)=c \tag{17}
\end{equation*}
$$

independent of $S$ and of the angle $\theta$ formed by the light propagation direction and the velocity of $S$ relative to $S_{0}$.

Consequences of A1. Consider the extreme points of the body in the directions $x$ and $y$ seen from $S_{0}$ at the same time $t_{0}$. From the first of (11) one gets:
$x_{1}=f_{1}\left(x_{01}-v t_{0}\right) ; x_{2}=f_{1}\left(x_{02}-v t_{0}\right)$
and from the second

$$
\begin{equation*}
y_{1}=g_{2} y_{01} ; \quad y_{2}=g_{2} y_{02} \tag{18}
\end{equation*}
$$

By subtracting from one another the two equations (18), and doing the same with (19), and comparing the so obtained results with (14) and (15) one obviously gets

$$
\begin{equation*}
f_{1}=\frac{1}{\sqrt{1-\beta^{2}}} \quad ; \quad g_{2}=1 \tag{20}
\end{equation*}
$$

Consequences of A2. Consider a clock at rest in some given point of $S$. Seen from $S_{0}$ the clock will obey the equation:

$$
x_{0}\left(t_{0}\right)=\beta c t_{0}+x_{0}(0)
$$

so that the fourth of (11) becomes

$$
\begin{equation*}
t=\left(e_{4}+e_{1} \beta c\right) t_{0}+e_{1} x_{0}(0) \tag{21}
\end{equation*}
$$

Considering any two times $t_{1}$ and $t_{2}$ marked by the moving clock and the corresponding $S_{0}$ times $t_{01}$ and $t_{02}$, one easily gets from (21):

$$
\begin{equation*}
t_{2}-t_{1}=\left(e_{4}+e_{1} \beta c\right)\left(t_{02}-t_{01}\right) \tag{22}
\end{equation*}
$$

Comparison with (16) gives

$$
\begin{equation*}
e_{4}+e_{1} \beta c=\sqrt{1-\beta^{2}} \tag{23}
\end{equation*}
$$

Consequences of A3. Obviously the angular dependence in (13) disappears only if

$$
\begin{equation*}
g_{2}=f_{1} \sqrt{1-\beta^{2}} \tag{24}
\end{equation*}
$$

after which (13) becomes

$$
\frac{1}{\tilde{c}_{2}(\theta)}=\frac{1}{c} \frac{e_{4}+e_{1} \beta c}{f_{1}\left(1-\beta^{2}\right)}
$$

Clearly $\tilde{c}_{2}(\theta)=c$ implies:

$$
\begin{equation*}
f_{1}\left(1-\beta^{2}\right)=e_{4}+e_{1} \beta c \tag{25}
\end{equation*}
$$

## 3. Equivalence of three pairs of assumptions

One can examine the consequences of A1, A2 and A3, by taking them two at a time in all possible ways.

Consequences of $\mathrm{A} 1+\mathrm{A} 2$. Considering together (20) and (23) one has:

$$
\begin{equation*}
f_{1}=\frac{1}{\sqrt{1-\beta^{2}}} ; \quad g_{2}=1 ; \quad e_{4}=\sqrt{1-\beta^{2}}-e_{1} \beta c \tag{26}
\end{equation*}
$$

Consequences of $\mathrm{A} 1+\mathrm{A} 3$. Considering together (20), (24) and (25) one has:

## 4. The inertial transformations

The meaning of (26), (27) and (28) is that the transformations of space and time relevant to the physical world are necessarily of the form:

$$
\left\{\begin{array}{l}
x=\frac{x_{0}-\beta c t_{0}}{\sqrt{1-\beta^{2}}}  \tag{29}\\
y=y_{0} \\
z=z_{0} \\
t=\sqrt{1-\beta^{2}} t_{0}+e_{1}\left(x_{0}-\beta c t_{0}\right)
\end{array}\right.
$$

where only one undetermined coefficient is left, $e_{1}$. As a consequence of (12) the inverse (one-way) velocity of light obtained for the values of the coefficients $f_{1}=\frac{1}{\sqrt{1-\beta^{2}}} ; \quad g_{2}=1 ; \quad e_{4}=\sqrt{1-\beta^{2}}-e_{1} \beta c$

Consequences of $\mathrm{A} 2+\mathrm{A} 3$. By inserting (23) in (25) one gets

$$
f_{1}=\frac{1}{\sqrt{1-\beta^{2}}}
$$

so that (24) gives

$$
g_{2}=1
$$

The last two conditions together with (23) are then

$$
\begin{equation*}
f_{1}=\frac{1}{\sqrt{1-\beta^{2}}} ; g_{2}=1 ; \quad e_{4}=\sqrt{1-\beta^{2}}-e_{1} \beta c \tag{28}
\end{equation*}
$$

The identity of (26), (27) and (28) constitutes the main point of the present paper: The three possible pairs of assumptions (A1, A2), (A1, A3) and (A2, A3) lead exactly to the same consequences. This is relevant to relativistic physics because (A1, A2) were the basic assumptions of Lorentz's reformulation of relativity [3]. Objections have been raised [4] against the validity of A1, for which there is indeed no direct experimental basis. There are, however, rather good experimental indications that A2 and A3 are true properties of nature [5-6]. Given the theorem just proved, the same indications can be taken as a rather convincing basis for the validity of A1 as well.

To this one can add that the Ehrenfest paradox [7], invoked as an argument against length contraction, is not a very serious problem in the real physical world. It is enough that the rotating disk becomes dome-shaped, in order to have a contracted circumference with a constant radius. Only in the abstract world of ideas, the symmetry of the problem between the two faces of the disk can be perfect. Their equivalence breaks down for a real disk, which is bound to have small irregularities. Furthermore, the reality of length contraction has been convincingly argued for with Bell's example of the thread connecting two equally accelerating spaceships [8].

$$
\begin{equation*}
\frac{1}{\tilde{c}_{1}(\theta)}=\frac{1}{c}+\left[\frac{\beta}{c}+e_{1} \sqrt{1-\beta^{2}}\right] \cos \theta \tag{27}
\end{equation*}
$$

The Lorentz transformations are recovered if one assumes $\tilde{c}_{1}(\theta)=c$. From (30) it follows:

$$
e_{1}=-\frac{1}{c} \frac{\beta}{\sqrt{1-\beta^{2}}}
$$

Different values of $e_{1}$ are obtained by using different clock synchronisation procedures. The so-called absolute synchronisation [9] is based on the idea that all clocks of $S$ are set to time $t=0$ when the passing clock at rest in the absolute system $S_{0}$ shows the time $t_{0}=0$. This means that from all positions in $S_{0}$ the time in $S$ will be seen to be the same, and therefore that no position dependent time-lag factor will be present in the transformation of time. Therefore $e_{1}=0$, condition which gives rise to a particularly simple transformation, different from the Lorentz one:

$$
\left\{\begin{array}{l}
x=\frac{x_{0}-\beta c t_{0}}{\sqrt{1-\beta^{2}}}  \tag{31}\\
y=y_{0} \\
z=z_{0} \\
t=\sqrt{1-\beta^{2}} t_{0}
\end{array}\right.
$$

The velocity of light relevant to a theory based on (31) is found by taking $e_{1}=0$ in (29):

$$
\frac{1}{\tilde{c}(\theta)}=\frac{1+\beta \cos \theta}{c}
$$

If (31) is inverted, it gives:

$$
\left\{\begin{array}{l}
x_{0}=\sqrt{1-\beta^{2}}\left[x+\frac{\beta c}{1-\beta^{2}} t\right]  \tag{33}\\
y_{0}=y \\
z_{0}=z \\
t_{0}=\frac{1}{\sqrt{1-\beta^{2}}} t
\end{array}\right.
$$

There is a formal difference between (31) and (33). The latter implies, for example, that the origin of $S_{0}$ (satisfying $x_{0}=y_{0}=z_{0}=0$ ) is described in $S$ by $y=z=0$ and by

$$
x=-\frac{\beta c}{1-\beta^{2}} t
$$

This origin is thus seen to move with speed $\beta c /\left(1-\beta^{2}\right)$ which can exceed $c$, but cannot be superluminal. In fact a light pulse seen from $S$ to propagate in the same direction as $S_{0}$ has $\theta=\pi$, and thus [using (32)] has velocity $\tilde{c}(\pi)=c /(1-\beta)$, which satisfies

$$
\frac{c}{1-\beta} \geq \frac{c \beta}{1-\beta^{2}}
$$

One of the typical features of these transformations is the presence of velocities which can grow without limit when they are relative to moving systems having absolute velocities $\beta c$ near to $c$. Absolute velocities can instead never exceed $c$ [10]. In STR one is used to relative velocities that are always equal and opposite, but this symmetry is a consequence of the particular synchronisation used and cannot be expected to hold more generally [10].

Consider now a third inertial system $S^{\prime}$ moving with velocity $\beta^{\prime} c$ and its transformation from $S_{0}$, which of course is

$$
\left\{\begin{array}{l}
x^{\prime}=\frac{x_{0}-\beta^{\prime} c t_{0}}{\sqrt{1-\beta^{\prime 2}}}  \tag{34}\\
y^{\prime}=y_{0} \\
z^{\prime}=z_{0} \\
t^{\prime}=\sqrt{1-\beta^{\prime 2}} t_{0}
\end{array}\right.
$$

By eliminating the $S_{0}$ variables from (34) and (33) one obtains the transformation between the two moving systems $S$ and $S^{\prime}$ :

$$
\left\{\begin{array}{l}
x^{\prime}=\frac{\sqrt{1-\beta^{2}}}{\sqrt{1-\beta^{\prime 2}}}\left[x-\frac{\left(\beta^{\prime}-\beta\right) c}{1-\beta^{2}} t\right]  \tag{35}\\
y^{\prime}=y \\
z^{\prime}=z \\
t^{\prime}=\frac{\sqrt{1-\beta^{\prime 2}}}{\sqrt{1-\beta^{2}}} t
\end{array}\right.
$$

The transformation (31) was first written down by Tangherlini [11]. For (33) and (35) see Ref. [10]. A possible name for (31)-(33)-(35) is "inertial transformations". In its most general form (35) an inertial transformation depends on two velocities ( $v$ and $v^{\prime}$ ). When one of them is zero, either $S$ or $S^{\prime}$ coincide with the privileged system $S_{0}$ and the transformation (35) becomes either (31) or (33). The inertial transformations have been shown not to form a group [10].

## 5. Some consequences of the inertial transformations

A feature characterising the transformations (31)-(33)-(35) is the existence of absolute simultaneity: two events taking place in different points of $S$ but at the same $t$ are judged to be simultaneous also in $S^{\prime}$ (and vice versa), this property being consequence of the absence of space variables in the transformation of time. Of course the existence of absolute simultaneity does not imply that time is absolute: in fact, the $\beta$-dependent factor in the transformation of time gives rise to time-dilation phenomena similar to those of STR. Time dilation in another sense is however also absolute: a clock at rest in $S$ is seen from $S_{0}$ to run slower, but a clock at rest in $S_{0}$ is seen from $S$ to run faster. Both observers agree that motion relative to $S_{0}$ slows down the pace of clocks, and the phenomenon loses the relativistic flavour it has in STR, becoming so to say absolute. Quantitatively one has for both situations:

$$
\begin{equation*}
\Delta t=\sqrt{1-\beta^{2}} \Delta t_{0} \tag{36}
\end{equation*}
$$

where $\Delta t$ and $\Delta t_{0}$ are the time intervals between any two given events as measured with clocks at rest in $S$ and in $S_{0}$, respectively. The difference with STR is however more apparent than real: a meaningful comparison of rates implies that a clock $T_{0}$ at rest in $S_{0}$ must be confronted with clocks at rest in different points of $S$. The result is thus dependent on the adopted convention for synchronising the latter clocks.

Absolute length contraction can also be deduced from (31)-(33). A rod at rest on the $x$ axis of $S$ between the points with co-ordinates $x_{2}$ and $x_{1}$ is seen in $S_{0}$ to have end points $x_{02}$ and $x_{01}$ at a common time $t_{0}$, where from (31):

$$
\begin{equation*}
x_{2}=\frac{x_{02}-v t_{0}}{\sqrt{1-\beta^{2}}} ; x_{1}=\frac{x_{01}-v t_{0}}{\sqrt{1-\beta^{2}}} \tag{37}
\end{equation*}
$$

From this one obtains

$$
\begin{equation*}
x_{2}-x_{1}=\frac{1}{\sqrt{1-\beta^{2}}}\left(x_{02}-x_{01}\right) \tag{38}
\end{equation*}
$$

The reasoning can be inverted by considering the rod at rest in $S$ and observed from $S_{0}$, and using (33). One gets, after a few simple steps:

$$
\begin{equation*}
x_{02}-x_{01}=\sqrt{1-\beta^{2}}\left(x_{2}-x_{1}\right) \tag{39}
\end{equation*}
$$

which could be obtained by inverting (38). The two results are thus mathematically equivalent and lead to the conclusion (with which both observers agree) that motion relative to $S_{0}$ leads to contraction. This is obviously an absolute effect, but again the discrepancy with the STR is due to the different conventions concerning clock synchronisation: the length of a moving rod can only be obtained by marking the simultaneous positions of its end points, and is therefore dependent on the very definition of simultaneity of distant events.

The assumed indifference of physical reality concerning synchronisation of clocks exists only insofar as one neglects accelerations: when these come into play every inertial system exists, so to say, only for a vanishingly small time interval and it is physically impossible to adopt any time-consuming procedure for the synchronisation of distant clocks in the accelerated frame (such as Einstein's procedure). Yet physical events take place and synchronisation is fixed by nature itself: the choice is $e_{1}=0$. How this happens was shown in Ref. [10].

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